

# Rellich estimates in $L^p$ for elliptic systems

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- joint work with Mihalis Mourgoglou, available on arXiv.
- earlier work with Andreas Axelsson, Alan McIntosh, Steve Hofmann

$\Omega = \mathbf{R}_+^{n+1}$ . Same analysis works in unit ball and every domain obtained by bilipschitz change of variables.

Points  $\mathbf{x} = (t, \mathbf{x})$ ,  $t > 0$ ,  $\mathbf{x} \in \mathbf{R}^n$ .

Measurable, bounded, with  $M_{m \times m}(\mathbf{C})$ -valued coefficients

$A_{i,j}$ ,  $i, j = 0, \dots, n$ ,  $m \geq 1$ .

Weak solution:  $u \in W_{loc}^{1,2}(\Omega; \mathbf{C}^m)$  and  $Lu = 0$  holds in  $\mathcal{D}'(\Omega; \mathbf{C}^m)$ :  
with summation convention

$$\operatorname{Re} \int_{\Omega} A_{i,j}^{\alpha,\beta} \partial_j u^\beta \overline{\partial_i \varphi^\alpha} = 0, \quad \forall \varphi \in C_0^\infty(\Omega; \mathbf{C}^m).$$

Short notation:  $A_{i,j}^{\alpha,\beta} \partial_j u^\beta \overline{\partial_i \varphi^\alpha} = A \nabla u \cdot \nabla \overline{\varphi}$  and  $Lu = \operatorname{div} A \nabla u$  in  $\Omega$ .

$i = 0$  corresponds to the vertical direction,  $i = 1, \dots, n$  to the horizontal directions.

- $\mathcal{E} := \dot{H}^1(\Omega; \mathbf{C}^m) =: \{u \in \mathcal{D}'(\Omega; \mathbf{C}^m) : \|\nabla u\|_2 < \infty\}$ .
- $\mathcal{E}/\mathbf{C}^m$  Banach space.
- $C_0^\infty(\bar{\Omega}; \mathbf{C}^m)$  dense in  $\mathcal{E}$  and  $\mathcal{E} \subset C([0, \infty); L_{\text{loc}}^2(\mathbf{R}^n; \mathbf{C}^m))$ .
- Trace of  $\mathcal{E}$  on  $\mathbf{R}^n$ :  $\mathcal{T} = \dot{H}^{1/2}(\mathbf{R}^n; \mathbf{C}^m)$ .  $\mathcal{T}$  contains  $C_0^\infty(\mathbf{R}^n)$ , dense.  $\mathcal{T} \subset L_{\text{loc}}^2(\mathbf{R}^n; \mathbf{C}^m)$ . Dual  $\dot{H}^{-1/2}(\mathbf{R}^n; \mathbf{C}^m)$ : space of distributions.
- $\mathcal{E}_0 = \{u \in \mathcal{E} : \text{Tr}(u) \in \mathbf{C}^m\}$ .  $C_0^\infty(\Omega; \mathbf{C}^m)$  dense in  $\mathcal{E}_0$ .
- When  $u \in \mathcal{E}$ ,  $Lu = 0$  (energy solution) means
$$\int_{\Omega} A \nabla u \cdot \nabla \bar{\varphi} = 0, \quad \forall \varphi \in \mathcal{E}_0.$$

If  $u \in \mathcal{E}$  and  $Lu = 0$ , then there exists a unique distribution  $g \in \mathcal{T}'$  such that

$$\int_{\Omega} A \nabla u \cdot \nabla \bar{\varphi} = \langle g, \varphi_0 \rangle, \quad \forall \varphi \in \mathcal{E}.$$

Notation:  $g = \partial_{\nu_A} u|_{t=0} = \partial_{\nu_A} u_0$ .

**Green's formula:** If  $u, w \in \mathcal{E}$  and  $Lu = 0 = L^*w$ , then

$$\langle u_0, \partial_{\nu_{A^*}} w_0 \rangle = \langle \partial_{\nu_A} u_0, w_0 \rangle. \quad (1)$$

Comparison between  $\|\nabla_{\tan} u_0\|$  and  $\|\partial_{\nu_A} u_0\|$  for energy solutions in some appropriate norm on the boundary.

Alternately, study of boundedness of the Dirichlet to Neumann operator or of the Neumann to Dirichlet operator.

**Ellipticity:** Assume for some  $\lambda > 0$ ,

$$\operatorname{Re} \int_{\Omega} A \nabla g \cdot \nabla \bar{g} \geq \lambda \int_{\Omega} |\nabla g|^2, \quad \forall g \in \mathcal{E}_0. \quad (2)$$

Let  $f \in \mathcal{T}$ . Then, there is a unique energy solution  $u \in \mathcal{E}$  of the equation  $\operatorname{div} A \nabla u = 0$  with  $u|_{t=0} = f$  (equality also in  $L^2_{\text{loc}}(\mathbf{R}^n; \mathbf{C}^m)$ , hence uniqueness). Moreover,

$$\|\nabla u\|_2 \sim \|f\|_{\mathcal{T}} \sim \|\nabla_{\tan} f\|_{\mathcal{T}'}$$

Hence, for any energy solution when  $A$  satisfies (2)

$$\|\partial_{\nu_A} u_0\|_{\mathcal{T}'} \lesssim \|\nabla_{\tan} u_0\|_{\mathcal{T}'}$$

notation:  $u$  has smooth Dirichlet datum if  $u_0 \in C_0^\infty$ .

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Let  $g \in \mathcal{T}'$ . Then, there is a energy solution  $u \in \mathcal{E}$ , unique modulo constants, of the equation  $\operatorname{div} A \nabla u = 0$  with  $\partial_{\nu_A} u|_{t=0} = g$ . Moreover,

$$\|\nabla u\|_2 \sim \|g\|_{\mathcal{T}'}$$

Hence, for any energy solution when  $A$  satisfies (3)

$$\|\nabla_{\tan} u_0\|_{\mathcal{T}'} \lesssim \|\partial_{\nu_A} u_0\|_{\mathcal{T}'}$$

Notation:  $u$  has smooth Neumann datum if  $g = \partial_{\nu_A} u_0 \in C_0^\infty$  with  $\int g = 0$ .



# BVP problems in $L^p$ , $1 < p < \infty$

Typical problems in harmonic analysis (for example for the Laplace equation).

• (Dir, A, p): Solve  $Lu = 0$  with  $\|\tilde{N}(u)\|_p < \infty$  and  $u_0 = f$  given in  $L^p(\mathbf{R}^n; \mathbf{C}^m)$ .

• (Reg, A, p): Solve  $Lu = 0$  with  $\|\tilde{N}(\nabla u)\|_p < \infty$  and  $\nabla_{\tan} u_0 = \nabla_{\tan} f$ ,  $f$  given in  $W^{1,p}(\mathbf{R}^n; \mathbf{C}^m)$ .

• (Neu, A, p): Solve  $Lu = 0$  with  $\|\tilde{N}(\nabla u)\|_p < \infty$  and  $\partial_{\nu_A} u|_{t=0} = g$  given in  $L^p(\mathbf{R}^n; \mathbf{C}^m)$ .

$\tilde{N}(h)$  is non-tangential maximal interior control of  $h$  defined in  $\Omega$ : it comes up quite naturally.

Not always solvable nor well-posed. No comprehensive theory at this time. To solve, Rellich inequalities are needed (not enough).

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Whitney ball:

$$W(t, x) := [(1 - c_0)t, (1 + c_0)t] \times B(x; c_1 t),$$

for fixed  $c_0 \in (0, 1)$ ,  $c_1 > 0$ .

$$\tilde{N}(h)(x) := \sup_{t>0} t^{-(n+1)/2} \|h\|_{L_2(W(t,x))}$$

## Theorem

Let  $A(\mathbf{x})$  be a bounded measurable matrix with the Gårding inequality (2). Let  $1 < p < \infty$ . The following are equivalent.

- 1 There exists  $C_p < \infty$  such that for any  $u \in \mathcal{E}$  solution of  $\operatorname{div} A \nabla u = 0$ ,  $\|\partial_{\nu_A} u_0\|_p \leq C_p \|\nabla_{\tan} u_0\|_p$ .
- 2 There exists  $C_{p'} < \infty$  such that for any  $w \in \mathcal{E}$  solution of  $\operatorname{div} A^* \nabla w = 0$ ,  $\|\partial_{\nu_{A^*}} w_0\|_{\dot{W}^{-1,p'}} \leq C_{p'} \|\nabla_{\tan} w_0\|_{\dot{W}^{-1,p'}}$ .

Moreover, in any direction it suffices the assumption holds for energy solutions with smooth Dirichlet data.

The tangential gradient and conormal derivative at the boundary are distributions in  $\mathbf{R}^n$  (in  $\mathcal{T}'$ ). Thus, finiteness of any of the norms above means that the distribution is identified with an element in the considered space which is also embedded in the space of distributions.

## Theorem

Let  $A(\mathbf{x})$  be a bounded measurable matrix with the stronger Gårding inequality (3). Let  $1 < p < \infty$ . The following are equivalent.

- 1 There exists  $C_p < \infty$  such that for any  $u \in \mathcal{E}$  solution of  $\operatorname{div} A \nabla u = 0$ ,  $\|\nabla_{\tan} u_0\|_p \leq C_p \|\partial_{\nu_A} u_0\|_p$ .
- 2 There exists  $C_{p'} < \infty$  such that for any  $w \in \mathcal{E}$  solution of  $\operatorname{div} A^* \nabla w = 0$ ,  $\|\nabla_{\tan} w_0\|_{\dot{W}^{-1,p'}} \leq C_{p'} \|\partial_{\nu_{A^*}} w_0\|_{\dot{W}^{-1,p'}}$ .

Moreover, in any direction it suffices the assumption holds for energy solutions with smooth Neumann data.



# Conclusion

There is a duality principle between Dirichlet problems in  $L^p$  and  $\dot{W}^{-1,p'}$  for  $L$  and  $L^*$  and a similar duality principle for Neumann problems in  $L^p$  and  $\dot{W}^{-1,p'}$  for  $L$  and  $L^*$ .

No assumption on  $A$  but bounded + elliptic. Systems OK.

One can formulate similar results for  $\frac{n}{n+1} < p \leq 1$  using Hardy and Hardy-Sobolev spaces and their duals.

**Problem:** Rellich not enough. Need also identification of the space of solutions corresponding to  $\|\nabla_{\tan} u_0\|_X + \|\partial_{\nu_A} u_0\|_X$ .

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## Theorem

(Auscher-Axelsson) Assume  $A(\mathbf{x}) = A(t, \mathbf{x}) = A(\mathbf{x})$  is independent of  $t$ , with the Garding inequality

$$\operatorname{Re} \int_{\mathbf{R}^n} A(\mathbf{x}) f(\mathbf{x}) \cdot \overline{f(\mathbf{x})} dx \geq \lambda \int_{\mathbf{R}^n} |f(\mathbf{x})|^2 dx, \quad \forall f \in \mathcal{H}_0, \quad (4)$$

where  $f \in \mathcal{H}_0$  means  $f = (f_{\perp}, f_{\parallel}) \in L^2(\mathbf{R}^n; \mathbf{C}^m \otimes (\mathbf{C}^m)^n)$  with  $\operatorname{curl}_{\tan} f_{\parallel} = 0$ . Then, for any weak solution, when the LHS is finite,

$$\|\tilde{N}(\nabla u)\|_2 \sim \|\nabla_{\tan} u_0\|_2 + \|\partial_{\nu_A} u_0\|_2.$$

$$\iint_{\Omega} t |\nabla u(t, \mathbf{x})|^2 dt dx \sim \|\nabla_{\tan} u_0\|_{\dot{W}^{-1,2}}^2 + \|\partial_{\nu_A} u_0\|_{\dot{W}^{-1,2}}^2.$$

Note: This holds also for  $A(t, \mathbf{x})$  satisfying (4) uniformly in  $t$  and a certain Carleson type condition measuring  $t$ -regularity.

## Theorem

(Jerison-Kenig and Kenig-Pipher for real equations, Auscher-Axelsson-McIntosh for systems) If  $A$  is  $t$ -independent with (4) and  $A = A^*$ , then for any weak solution with  $\|\tilde{N}(\nabla u)\|_2 < \infty$  or any energy solution

$$\|\nabla_{\tan} u_0\|_2 \sim \|\partial_{\nu_A} u_0\|_2.$$

Furthermore,  $(Dir', A, 2)$ ,  $(Reg, A, 2)$ ,  $(Neu, A, 2)$  are well-posed.

Here,  $(Dir', A, 2)$  is formulated as follows:  $Lu = 0$ ,  $u_0 = f$  given in  $L^2(\mathbf{R}^n; \mathbf{C}^m)$  with interior control  $\iint_{\Omega} t|\nabla u(t, x)|^2 dt dx < \infty$  instead of  $\|\tilde{N}(u)\|_2 < \infty$ .

This comparison follows from an integral identity in the spirit of the identities discovered by Rellich for eigenvalue problems. Self-adjointness is essential.

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This comparison follows from an integral identity in the spirit of the identities discovered by Rellich for eigenvalue problems. Self-adjointness is essential.

## Theorem

(Auscher-McIntosh-Mourgoglou) If  $A$  is  $t$ -independent with (4) and  $A$  **block lower-triangular** ( $A_{0,j} = 0$  for  $j = 1, \dots, n$ ), then for any weak solution with  $\|\tilde{N}(\nabla u)\|_2 < \infty$  or any energy solution

$$\|\nabla_{\tan} u_0\|_2 \lesssim \|\partial_{\nu_A} u_0\|_2.$$

Furthermore,  $(\text{Neu}, A, 2)$  is well-posed.

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$$\|\partial_{\nu_A} u_0\|_2 \lesssim \|\nabla_{\tan} u_0\|_2.$$

Furthermore,  $(\text{Reg}, A, 2)$  and  $(\text{Dir}', A^*, 2)$  are well-posed.



## Theorem

(Hofmann-Kenig-Mayboroda-Pipher, 2012 & 2013) If  $A$  is  $t$ -independent, **real with  $m = 1$  (equation)**, then there is a  $2 \leq p_0 < \infty$  such that  $(\text{Dir}, A, p)$  is well-posed and  $(\text{Reg}, A^*, p')$  is well-posed when  $p_0 < p < \infty$ .

To prove Dirichlet, they showed for any energy solution

$$\|\partial_{\nu_A} u_0\|_{\dot{W}^{-1,p}} \lesssim \|\nabla_{\tan} u_0\|_{\dot{W}^{-1,p}}. \quad (5)$$

Their derivation of  $(\text{Reg}, A^*, p')$  is quite technical. Here is a simple (at least conceptually) way for existence: We have (A.-Mourgoglou) for any energy solution of  $L^* u = 0$  and  $1 < p' \leq 2$ ,

$$\|\tilde{N}(\nabla u)\|_{p'} \sim \|\nabla_{\tan} u_0\|_{p'} + \|\partial_{\nu_{A^*}} u_0\|_{p'}.$$

Apply (5) and our duality result:

$$\|\partial_{\nu_{A^*}} u_0\|_{p'} \lesssim \|\nabla_{\tan} u_0\|_{p'}.$$

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